# Quantum Mechanics

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**Sheet 11 Saturday December 21 at 24:00**

#### **1** *p***-orbitals**

In this exercise we investigate the spherical harmonics corresponding to the angular momentum quantum number  $l = 1$ . However, let us first briefly consider the case  $l = 0$ . Since  $-l \le m \le l$ , the only possibility is  $m = 0$ , and the corresponding spherical harmonics is  $Y_{l=0}^{m=0}(\theta, \phi) = (4\pi)^{-1/2}$ . One can think of  $Y_0^0$  as a uniform sphere, and is commonly referred to as an *s*-orbital. The next possibility is  $l = 1$ , for which  $m = -1, 0, 1$ , and which correspond to the *p*-orbitals.

**a)** In spherical coordinates the orbital angular momentum operators take the form

$$
\langle \theta, \phi | L_x | \psi \rangle = i\hbar \left( \sin \phi \frac{\partial}{\partial \theta} + \frac{\cos \phi}{\tan \theta} \frac{\partial}{\partial \phi} \right) \psi(\theta, \phi),
$$
  

$$
\langle \theta, \phi | L_y | \psi \rangle = i\hbar \left( -\cos \phi \frac{\partial}{\partial \theta} + \frac{\sin \phi}{\tan \theta} \frac{\partial}{\partial \phi} \right) \psi(\theta, \phi),
$$
  

$$
\langle \theta, \phi | L_z | \psi \rangle = -i\hbar \frac{\partial}{\partial \phi} \psi(\theta, \phi).
$$

*Show that the ladder operators L*±*, expressed in the spherical coordinate system, are*

$$
\langle \theta, \phi | L_{\pm} | \psi \rangle = \hbar e^{\pm i \phi} \Big( \pm \frac{\partial}{\partial \theta} + i \frac{1}{\tan \theta} \frac{\partial}{\partial \phi} \Big) \psi(\theta, \phi).
$$

**(2 points)**

**b)** It is the case that

$$
Y_1^1(\theta,\phi)=-\sqrt{\frac{3}{8\pi}}\sin\theta e^{i\phi},
$$

where one might note that  $Y_1^1(\theta, \phi) = \langle \theta, \phi | l = 1, m = 1 \rangle$ . Apply the ladder operators onto  $Y_1^1$ , in *order to show that*

$$
Y_1^0(\theta,\phi) = \sqrt{\frac{3}{4\pi}} \cos \theta, \quad Y_1^{-1}(\theta,\phi) = \sqrt{\frac{3}{8\pi}} \sin \theta e^{-i\phi}.
$$

#### **(4 points)**

- <span id="page-0-0"></span>**c)** In the following we take a closer look at these functions, and construct the corresponding *p*-orbitals. As the first step, we put  $p_z(\theta, \phi) = Y_1^0(\theta, \phi)$ . *Sketch the graph of*  $|p_z|$  *in the x-z-plane, i.e., make a polar plot of*  $|p_z(\theta, \phi = 0)|$ . (Hence, at the angle  $\theta$  with respect to the *z*-axis, put a point at the radius |*pz*(*θ*, 0)|.) *Next, sketch the shape of* |*pz*| *in* **R**<sup>3</sup> *.* Rough sketches are enough (but you can use some plotting tool if you want to). **(2 points)**
- **d**) The complex valued functions  $Y_1^{\pm 1}$  are a bit harder to interpret, but we can make linear combinations of them in order to get real-valued functions. *Determine the functions*

$$
p_x(\theta, \phi) = -\frac{1}{\sqrt{2}} (Y_1^1(\theta, \phi) - Y_1^{-1}(\theta, \phi)), \quad p_y(\theta, \phi) = -\frac{1}{i\sqrt{2}} (Y_1^1(\theta, \phi) + Y_1^{-1}(\theta, \phi)).
$$

**(2 points)**

#### **(2 points)**

**Remark:** Since  $p_x$ ,  $p_y$ ,  $p_z$  and  $Y_1^1$ ,  $Y_1^0$ ,  $Y_1^{-1}$  both span  $\mathcal{H}_l$ , it is a matter of convenience which one to use. It turns out that the *p*-orbitals  $p_x$ ,  $p_y$ ,  $p_z$  are convenient for molecular physics and chemistry. In particular, these orbitals distribute the particle<sup>[1](#page-1-0)</sup> (the electron, in case of chemistry) in distinct manners over space, which is useful when describing the formation of chemical bonds. In  $\mathbf{c}$  you have already seen how  $p_x$  is distributed, and in the next problem we consider  $p_y$  and  $p_z$ .

**f)** Sketch the graph of  $|p_x|$  in the x-z-plane, i.e., make a polar plot of  $|p_x(\theta, \phi = 0)|$ . Similarly, sketch the *graph of*  $|p_y|$  *in the y-z-plane, i.e., make a polar plot of*  $|p_y(\theta, \phi = \pi/2)|$ *.* 

# **(2 points)**

**Remark:** In case you are using some plot-tool, it is instructive to make 3D plots of  $|p_x|$  and  $|p_y|$  and compare with your plot of  $|p_z|$ .

**g**) The orbital  $p_z$  is an eigenfunction of  $L_z$ . However, neither  $p_x$  nor  $p_y$  are eigenfunctions of  $L_z$ (since both are linear combinations of eigenfunctions with different eigenvalues). *Show that p<sup>x</sup> is an eigenfunction of L<sup>x</sup> and that p<sup>y</sup> is an eigenfunction of Ly. What are the eigenvalues?*

# **(4 points)**

**h)** *Show that p<sup>x</sup> is symmetric under rotations around the x-axis and show that p<sup>y</sup> is symmetric under rotations around the y-axis.* Looking at the concrete functions  $p_x$  and  $p_y$ , these symmetries are not obvious (at least not to me). However, if you think about what you have learned, there is a very direct argument.

### **(2 points)**

<span id="page-1-0"></span><sup>&</sup>lt;sup>1</sup>To be more precise, the *p*'s and the *Y*'s only give the angular part of the wave function, and are accompanied by radial wave-functions. The absolute value square of the joint wave-function gives the probability density of where to find the electron. For the p-orbitals, these distributions are oriented in a distinct manner around the atom. This can be compared with the *s*-orbital, where the probability distribution is rotation symmetric.