QUANTUM MECHANICS

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1 *p*-orbitals

In this exercise we investigate the spherical harmonics corresponding to the angular momentum quantum number l = 1. However, let us first briefly consider the case l = 0. Since $-l \le m \le l$, the only possibility is m = 0, and the corresponding spherical harmonics is $Y_{l=0}^{m=0}(\theta, \phi) = (4\pi)^{-1/2}$. One can think of Y_0^0 as a uniform sphere, and is commonly referred to as an *s*-orbital. The next possibility is l = 1, for which m = -1, 0, 1, and which correspond to the *p*-orbitals.

a) In spherical coordinates the orbital angular momentum operators take the form

$$\begin{split} \langle \theta, \phi | L_x | \psi \rangle &= i\hbar \bigg(\sin \phi \frac{\partial}{\partial \theta} + \frac{\cos \phi}{\tan \theta} \frac{\partial}{\partial \phi} \bigg) \psi(\theta, \phi), \\ \langle \theta, \phi | L_y | \psi \rangle &= i\hbar \bigg(-\cos \phi \frac{\partial}{\partial \theta} + \frac{\sin \phi}{\tan \theta} \frac{\partial}{\partial \phi} \bigg) \psi(\theta, \phi), \\ \langle \theta, \phi | L_z | \psi \rangle &= -i\hbar \frac{\partial}{\partial \phi} \psi(\theta, \phi). \end{split}$$

Show that the ladder operators L_{\pm} , expressed in the spherical coordinate system, are

$$\langle \theta, \phi | L_{\pm} | \psi \rangle = \hbar e^{\pm i\phi} \Big(\pm \frac{\partial}{\partial \theta} + i \frac{1}{\tan \theta} \frac{\partial}{\partial \phi} \Big) \psi(\theta, \phi)$$

(2 points)

b) It is the case that

$$Y_1^1(heta,\phi) = -\sqrt{rac{3}{8\pi}}\sin heta e^{i\phi},$$

where one might note that $Y_1^1(\theta, \phi) = \langle \theta, \phi | l = 1, m = 1 \rangle$. Apply the ladder operators onto Y_1^1 , in order to show that

$$Y_1^0(\theta,\phi) = \sqrt{\frac{3}{4\pi}}\cos\theta, \quad Y_1^{-1}(\theta,\phi) = \sqrt{\frac{3}{8\pi}}\sin\theta e^{-i\phi}.$$

(4 points)

- c) In the following we take a closer look at these functions, and construct the corresponding *p*-orbitals. As the first step, we put p_z(θ, φ) = Y₁⁰(θ, φ). Sketch the graph of |p_z| in the x-z-plane, *i.e., make a polar plot of* |p_z(θ, φ = 0)|. (Hence, at the angle θ with respect to the z-axis, put a point at the radius |p_z(θ, 0)|.) *Next, sketch the shape of* |p_z| in ℝ³. Rough sketches are enough (but you can use some plotting tool if you want to).
- **d)** The complex valued functions $Y_1^{\pm 1}$ are a bit harder to interpret, but we can make linear combinations of them in order to get real-valued functions. *Determine the functions*

$$p_{x}(\theta,\phi) = -\frac{1}{\sqrt{2}} \left(Y_{1}^{1}(\theta,\phi) - Y_{1}^{-1}(\theta,\phi) \right), \quad p_{y}(\theta,\phi) = -\frac{1}{i\sqrt{2}} \left(Y_{1}^{1}(\theta,\phi) + Y_{1}^{-1}(\theta,\phi) \right).$$

(2 points)

(2 points)

Remark: Since p_x , p_y , p_z and Y_1^1 , Y_1^0 , Y_1^{-1} both span \mathcal{H}_l , it is a matter of convenience which one to use. It turns out that the *p*-orbitals p_x , p_y , p_z are convenient for molecular physics and chemistry. In particular, these orbitals distribute the particle¹ (the electron, in case of chemistry) in distinct manners over space, which is useful when describing the formation of chemical bonds. In **c**) you have already seen how p_x is distributed, and in the next problem we consider p_y and p_z .

f) Sketch the graph of $|p_x|$ in the x-z-plane, i.e., make a polar plot of $|p_x(\theta, \phi = 0)|$. Similarly, sketch the graph of $|p_y|$ in the y-z-plane, i.e., make a polar plot of $|p_y(\theta, \phi = \pi/2)|$.

(2 points)

Remark: In case you are using some plot-tool, it is instructive to make 3D plots of $|p_x|$ and $|p_y|$ and compare with your plot of $|p_z|$.

g) The orbital p_z is an eigenfunction of L_z . However, neither p_x nor p_y are eigenfunctions of L_z (since both are linear combinations of eigenfunctions with different eigenvalues). Show that p_x is an eigenfunction of L_x and that p_y is an eigenfunction of L_y . What are the eigenvalues?

(4 points)

h) Show that p_x is symmetric under rotations around the *x*-axis and show that p_y is symmetric under rotations around the *y*-axis. Looking at the concrete functions p_x and p_y , these symmetries are not obvious (at least not to me). However, if you think about what you have learned, there is a very direct argument.

(2 points)

¹To be more precise, the *p*'s and the *Y*'s only give the angular part of the wave function, and are accompanied by radial wave-functions. The absolute value square of the joint wave-function gives the probability density of where to find the electron. For the p-orbitals, these distributions are oriented in a distinct manner around the atom. This can be compared with the *s*-orbital, where the probability distribution is rotation symmetric.